



Bound states & scattering states



- Bound states & scattering states
- Real potentials



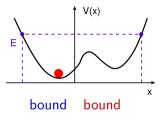
- Bound states & scattering states
- Real potentials
- The Dirac delta function



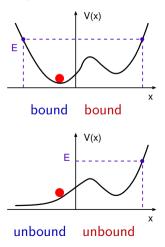
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- The delta function well



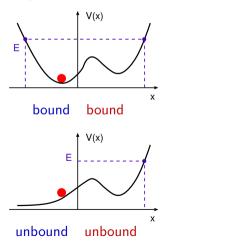


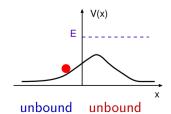




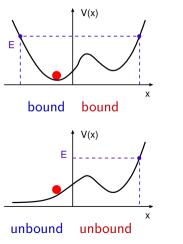


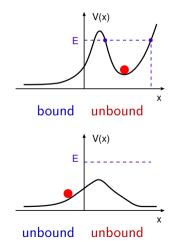












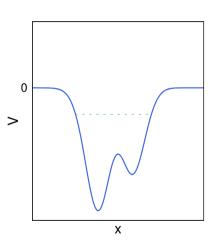


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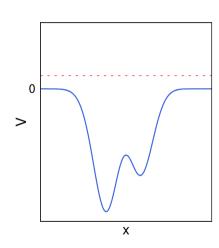




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When E < 0 we will have bound states with discrete energy levels if the potential is negative anywhere in space

When E > 0 we will have an unbound system with continuous energies and scattering

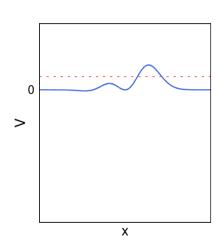




Real potentials will always trend toward zero at large values of x and so we have a much simpler situation

When E < 0 we will have bound states with discrete energy levels if the potential is negative anywhere in space

When E > 0 we will have an unbound system with continuous energies and scattering (independent of the sign of the potential!)





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When multiplied with a function it picks out the function at a single value

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This integral works for any limits which include the peak of the delta function.



Consider a potential of the form



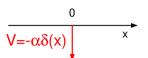
Consider a potential of the form

$$\frac{0}{V=-\alpha\delta(x)} \downarrow x$$

$$V(x) = -\alpha \delta(x)$$



Consider a potential of the form where α is a positive constant



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Consider a potential of the form where α is a positive constant

the Schrödinger equation is now and if E < 0, there is a bound state

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Consider a potential of the form where α is a positive constant

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$$\frac{d^2\psi}{dx^2} = -\frac{2mE}{\hbar^2}\psi = \kappa^2\psi$$



Consider a potential of the form where α is a positive constant

the Schrödinger equation is now and if E < 0, there is a bound state start with region x < 0 where the potential is zero and $\kappa > 0$

$$V(x) = -\alpha \delta(x)$$

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 $E\psi = -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} - \alpha\delta(x)\psi$ $\frac{d^2\psi}{dx^2} = -\frac{2mE}{\hbar^2}\psi = \kappa^2\psi$ $\kappa \equiv \frac{\sqrt{-2mE}}{\hbar} > 0$

the general solution is



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$$V(x) = -\alpha \delta(x)$$

$$E\psi = -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} - \alpha\delta(x)\psi$$
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$$\kappa \equiv \frac{\sqrt{-2mE}}{\hbar} > 0$$
$$\psi(x) = Ae^{-\kappa x} + Be^{+\kappa x}$$

Delta function potential well



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$$= Be^{+\kappa x}, \quad (x < 0)$$



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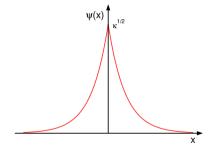
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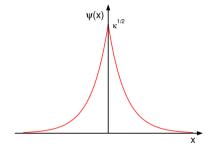
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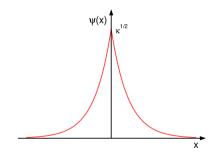
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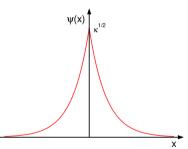
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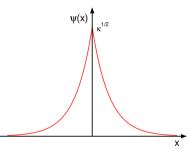
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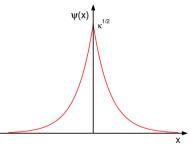
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$$E\int_{-\epsilon}^{+\epsilon} \psi(x)dx = -\frac{\hbar^2}{2m}\int_{-\epsilon}^{+\epsilon} \frac{d^2\psi(x)}{dx^2}dx + \int_{-\epsilon}^{+\epsilon} V(x)\psi(x)dx$$



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The last term is usually zero, unless $V(x) \to \infty$



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For the delta function, this limit is non-zero and can be calculated

For our solution

$$\psi(x) = \begin{cases} Be^{-\kappa x}, & x \ge 0 \\ Be^{+\kappa x}, & x \le 0 \end{cases}$$

$$\Delta \left(\frac{d\psi}{dx}\right) = \frac{2m}{\hbar^2} \lim_{\epsilon \to 0} \int_{-\epsilon}^{+\epsilon} \frac{V(x)\psi(x)dx}{V(x)\psi(x)dx}$$
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$$\frac{d\psi}{dx} = \begin{cases} -B\kappa e^{-\kappa x}, & x > 0\\ +B\kappa e^{+\kappa x}, & x < 0 \end{cases}$$

$$\Delta \left(\frac{d\psi}{dx}\right) = \frac{2m}{\hbar^2} \lim_{\epsilon \to 0} \int_{-\epsilon}^{+\epsilon} \frac{V(x)\psi(x)dx}{V(x)\psi(x)dx}$$
$$= \frac{2m}{\hbar^2} \lim_{\epsilon \to 0} \int_{-\epsilon}^{+\epsilon} \frac{-\alpha\delta(x)\psi(x)dx}{\delta(x)\psi(x)dx}$$
$$= -\frac{2m\alpha}{\hbar^2} \psi(0) \lim_{\epsilon \to 0} \int_{-\epsilon}^{+\epsilon} \frac{\delta(x)dx}{\delta(x)dx}$$
$$= -\frac{2m\alpha}{\hbar^2} \psi(0)$$



For the delta function, this limit is non-zero and can be calculated

For our solution

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Delta function discontinuity



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$$\kappa = \frac{m\alpha}{\hbar^2} \rightarrow E = -\frac{\hbar^2 \kappa^2}{2m} = -\frac{m\alpha^2}{2\hbar^2}$$



Thus, the negative delta-function potential has a single bound state with wave function



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Thus, the negative delta-function potential has a single bound state with wave function and energy.

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Thus, the negative delta-function potential has a single bound state with wave function and energy. There is always only one bound state for this potential, independent of the strength of the potential α .

$$\psi(x) = \frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2}; \quad E = -\frac{m\alpha^2}{2\hbar^2}$$





General solution for three regions



- General solution for three regions
- Applying the boundary conditions

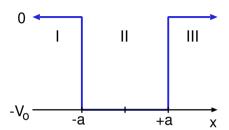


- General solution for three regions
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- Even solutions

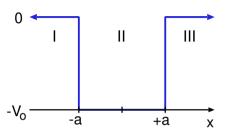


- General solution for three regions
- Applying the boundary conditions
- Even solutions
- Limiting cases



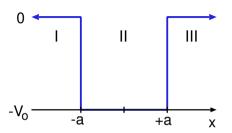






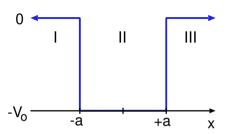
In region I, x < -a





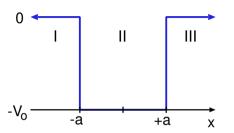
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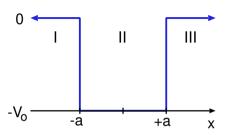


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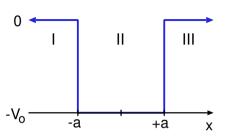


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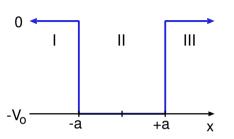
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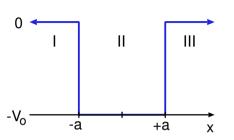
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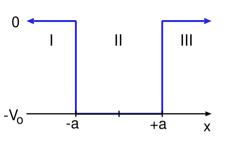
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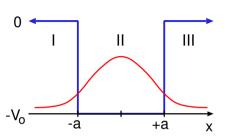
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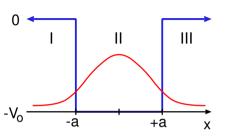
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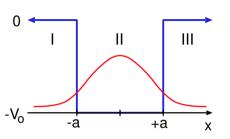
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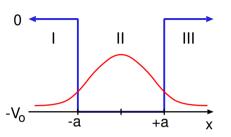
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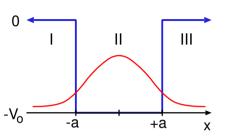
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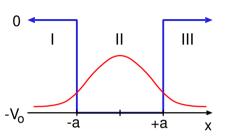
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Both the wave function and its derivative must be continuous at the boundaries of the three regions, x = -a, +a.





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Both the wave function and its derivative must be continuous at the boundaries of the three regions, x = -a, +a. Let's consider the even solutions initially, where $C \equiv 0$.

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$$\kappa^{2} + I^{2} = \frac{-2mE}{\hbar^{2}} + \frac{2m(E + V_{0})}{\hbar^{2}}$$
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dividing the two equations:

at $x = \pm a$ we have

$$\kappa^2 + I^2 = \frac{-2mE}{\hbar^2} + \frac{2m(E + V_0)}{\hbar^2}$$

$$\kappa^2 + \frac{z^2}{r^2} = \frac{2mV_0}{\hbar^2} = \frac{z_0^2}{r^2}$$

$$\kappa = I \tan(Ia)$$



Both the wave function and its derivative must be continuous at the boundaries of the three regions, x = -a, +a. Let's consider the even solutions initially, where $C \equiv 0$.

$$\psi(x) = \begin{cases} Be^{+\kappa x}, & x < -a \\ D\cos(Ix), & |x| < a \\ Fe^{-\kappa x}, & x > +a \end{cases}$$

at
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 we have

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at x = +a we have

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$$\frac{1}{z} \sqrt{z_0^2 - z^2} = \tan z$$

$$\sqrt{\left(\frac{z_0}{z}\right)^2 - 1} = \tan z$$

Even solutions to finite well



$$\tan z = \sqrt{\left(\frac{z_0}{z}\right)^2 - 1}$$

Even solutions to finite well



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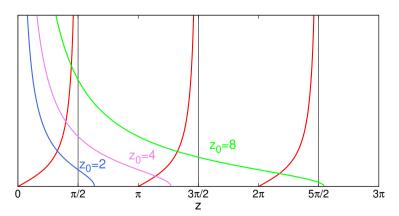
This is a transcendental equation which defines the discrete energies which are allowed as stationary states.

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and the even solutions approach

$$z = Ia = \frac{\sqrt{2m(E + V_0)}}{\hbar}a$$

Since $E_n + V_0$ is just the energy above the bottom of the well and the width is 2a, the even solutions (and the odd ones, of course) approach those of the infinite square well.

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 then $z_0 o \infty$
$$z_n o \frac{n\pi}{2}, \quad n=1,3,5,\cdots$$

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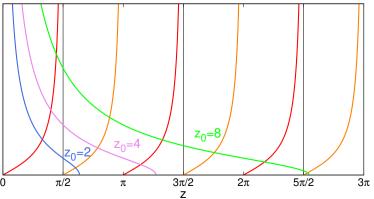


As the well becomes more shallow, $V_0 \rightarrow 0$ and $z_0 \rightarrow 0$ as well.

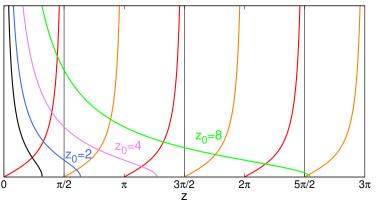


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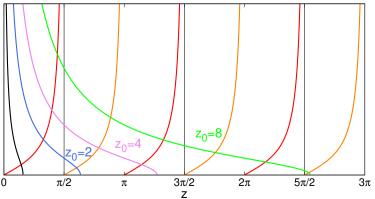
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The number of states decreases until the lowest odd bound state vanishes. However, the ground state (lowest even state) will never vanish. There is *always* a bound state no matter how shallow the well!

